

# PHYS 705: Classical Mechanics



## Pre-class notes:

- Double Pendulum demo
- Catch up week: you will have two weeks to do HW#5 (due Oct 12)
- Due to Columbus Day Holiday, class will meet on Tuesday (Oct 12) instead of Monday (Oct 11)

## HW#3:

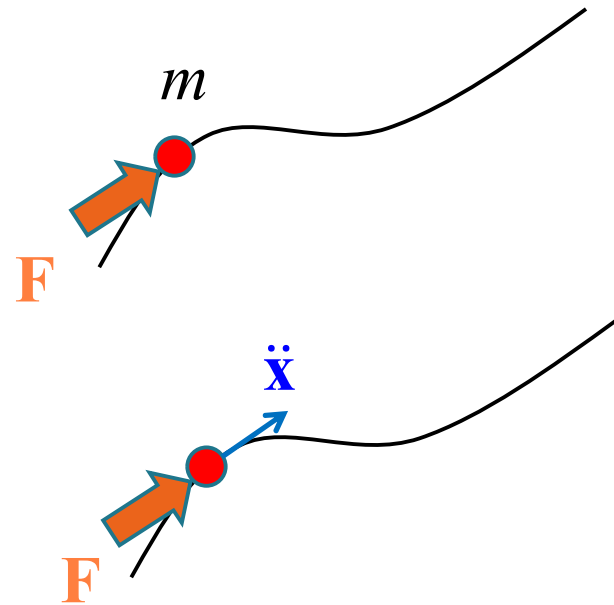
Newton's 2<sup>nd</sup> Law:

- Mass  $m$  is being pushed by a net force  $\mathbf{F}$
- Its motion (its acceleration)  $\ddot{\mathbf{x}}$  is given by the 2<sup>nd</sup> law

$$\frac{d\mathbf{P}}{dt} = m\ddot{\mathbf{x}} = \mathbf{F}$$

how it moves

the push



**$m\ddot{\mathbf{x}}$  is NOT a force !**

## Conservation and Symmetry

There is a very strong link between Symmetry and Conservation Theorems:

- Symmetry  conserved quantity
- New Terminology:

$$\left. \begin{array}{l} \text{generalized momentum} \\ \text{canonical momentum} \\ \text{conjugate momentum} \end{array} \right\} \equiv \boxed{p_j \equiv \frac{\partial L}{\partial \dot{q}_j}}$$

$$\text{cyclic coordinate} \equiv \text{one that doesn't appear in } L, \text{ i.e., } \frac{\partial L}{\partial q_j} = 0$$

## Conservation and Symmetry

Let consider a system in which all applied forces are derivable from a scalar potential function and we have chosen a set of “proper” generalized coordinates (i.e., all constraints are “hidden”). Then,

$$\frac{\partial L}{\partial q_j} - \frac{d}{dt} \left( \frac{\partial L}{\partial \dot{q}_j} \right) = 0$$

A generalized coordinate  $q_j$  is **cyclic** if  $\frac{\partial L}{\partial q_j} = 0$

Then, EL equation implies  $\frac{d}{dt} \left( \frac{\partial L}{\partial \dot{q}_j} \right) = 0$

$$\longrightarrow p_j \equiv \frac{\partial L}{\partial \dot{q}_j} = \text{const} \quad p_j \text{ is conserved !}$$

# Conservation and Symmetry

Simple examples:

→ If  $q_j$  corresponds to a **translation** relative to fixed Cartesian coordinates



Conservation of Linear Momentum

→ If  $q_j$  corresponds to a **rotation** of the system



Conservation of Angular Momentum

→ In general, if  $q_j$  is a cyclic variable, its corresponding conjugate momentum  $p_j$  will be conserved.

# Energy Conservation and Time Invariance

Defining  $h = \sum_j \dot{q}_j \frac{\partial L}{\partial \dot{q}_j} - L$

1. If  $t$  is a cyclic variable ( $L$  is **time invariance**),  $h$  is conserved.

$$\frac{\partial L}{\partial t} = 0 \quad \Rightarrow \quad \frac{dh}{dt} = -\frac{\partial L}{\partial t} = 0 \quad h \text{ is constant in time!}$$

2. Additionally, if

$$U = U(q_i) \text{ (} U \text{ does not dep on } \dot{q}_i \text{) AND}$$

$$\frac{\partial \mathbf{r}_i}{\partial t} = 0 \text{ (the coord trans does not dep on time explicitly)}$$

$$\Rightarrow h = E \text{ (total energy)!}$$

## Energy Conservation and Time Invariance

$$h(q_j, \dot{q}_j, t) = \sum_j \dot{q}_j \frac{\partial L}{\partial \dot{q}_j} - L$$

is call the **Jacobi Integral** or the energy function.

$h$  is closely related to the **Hamiltonian function**  $H(q_i, p_i, t)$  that we will introduce later in the semester.



$h(q_i, \dot{q}_i, t)$  is a function of the generalized coordinates and their time derivatives

But,  $H(q_i, p_i, t)$  is a function of the generalized coordinates and the associated generalized momenta